THE GEOMETRICAL MULTISCALE MODELING IN HEMODYNAMICS

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Workshop on Nonlinear PDE's

Dedicated to 80th anniversary of birth of Prof. Jindřich Nečas

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MOTIVATION

Hemodynamics vs cardiovascular diseases: local fluid patterns and **wall shear stress** are strictly related to the development of cardiovascular diseases (indicator of atherosclerosis)

Difficulties in modeling blood flow

- Blood Rheology
- Complex geometry



3D flow simulations are restricted to specific regions of interest

Closed system

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Local flow dynamics has an important role in the systemic circulation (and vice-versa)



WSS pulmonary artery (congenital heart disease)



MATHEMATICAL MODEL

- 1. Blood rheology
- 2. Fluid-structure interaction
- 3. Geometrical multiscale approach





CIRCULATORY SYSTEM: FLUID DYNAMIC VALUES

Relationship between arterial size, number of vessels and average Reynolds numbers

Vessel	Radius (cm)	Number	Wall thickness (cm)	Average Re number
Aorta	1.25	1	0.2	3400
Arteries	0.2	159	0.1	500
Arterioles	1.5×10⁻³	5.7×10 ⁷	2×10⁻³	0.7
Capillaries	3×10⁻⁴	1.6×10 ¹⁰	1×10 ⁻⁴	0.002
Venules	1×10 ⁻³	1.3×10 ⁹	2×10 ⁻⁴	0.01
Veins	0.25	200	0.05	140
Vena cava	1.5	1	0.15	3300

Turbulence can develop in a few cases:

High cardiac output (exercise); Stenoses; Low blood density (for example: anemia)





BLOOD COMPOSITION

- Blood is a suspension of
 - cells
 - erythrocytes (RBCs)
 - leukocytes (WBCs)
 - platelets
 - plasma (90-92% water + proteins, organic salts)



Plasma Red blood cells Platelets White blood cells

	Number/ mm ³	Shape (unstressed)	Size μm (unstressed)	Volume Conc.
erythrocytes	4-6×10 ⁶	Biconcave discs with no nuclei	8×1-3	45%
leukocytes	4-11 ×10 ³	roughly spherical	7-22	1%
platelets	2.5-5 ×10⁵	Rounded or oval discs	2-4	





BLOOD RHEOLOGY

Why is blood a non-Newtonian fluid ?

Non-Constant Viscosity



NEWTONIAN vs NON-NEWTONIAN FLUID BEHAVIOR

Non - Constant Viscosity

$$T = -pI + \tau = -pI + 2\mu(\dot{\gamma})D$$

- Shear thinning (or pseudoplastic) fluids
- Shear thickening (or dilatant) fluids
- Yield stress (Bingham plastic) fluids



cena



$$\mu(\gamma) = \frac{\tau(\gamma)}{\gamma}, \qquad \mu(\gamma) > 0$$

shear-viscosity function (apparent viscosity)

at constant shear rate:

- Thixotropic fluids (apparent viscosity decreasing in time)
- Rheopectic fluids (apparent viscosity increasing in time)



SHEAR-THINNING BLOOD FLOW MODELS

$\mu_0 = \lim_{\dot{\gamma} o 0} \mu(\dot{\gamma}) =$	= 0.056 <i>Pas</i>	$\mu_{\infty} = \lim_{\dot{\gamma} o \infty} \mu(\dot{\gamma}) = 0.00345 Pas$
Model	$rac{\mu(\dot{\gamma})-\mu_\infty}{\mu_0-\mu_\infty}$	MATERIAL CONSTANTS FOR BLOOD
Powell-Eyring	$\frac{\sinh^{-1}(\lambda\dot{\gamma})}{\lambda\dot{\gamma}}$	$\lambda = 5.383s$
Cross	$(1+(\lambda\dot\gamma)^m)^{-1}$	$\lambda = 1.007s, m = 1.028$
Modified Cross	$(1+(\lambda\dot\gamma)^m)^{-a}$	$\lambda = 3.736s, m = 2.406, a = 0.254$
Carreau	$(1+(\lambda\dot{\gamma})^2)^{(n-1)/2}$	$\lambda=3.313s,n=0.3568$
CARREAU-YASUDA	$(1+(\lambda\dot\gamma)^a)^{(n-1)/a}$	$\lambda = 1.902s, n = 0.22, a = 1.25$

(Y.I.Cho and K.R.Kensey, *Biorheology*, 1991)

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BLOOD RHEOLOGY

- Anne M. Robertson, Adélia Sequeira and Marina V. Kameneva. Hemorheology. In: *Hemodynamical Flows: Modeling, Analysis and Simulation*, G. P. Galdi, R. Rannacher, A. M. Robertson, S. Turek, Oberwolfach Seminars, Vol. 37, pp.63-120, 2008.
- Anne M. Robertson, Adélia Sequeira and Robert Owens. Rheological models for blood. In: Cardiovascular Mathematics, A. Quarteroni, L. Formaggia and A. Veneziani (eds.), Springer-Verlag, 2009.





BLOOD FLOW DYNAMICS

Blood flow: Generalized Newtonian fluid equations

$$\rho \left(\frac{\partial \boldsymbol{u}}{\partial t} + \boldsymbol{u} \cdot \nabla \boldsymbol{u} \right) + \nabla p - \nabla \cdot \boldsymbol{\tau}(\boldsymbol{u}) = 0 \quad \text{in } \Omega$$
$$\nabla \cdot \boldsymbol{u} = 0 \quad \text{in } \Omega$$

$$\boldsymbol{\tau}(\boldsymbol{u}) = 2\mu(\dot{\gamma})\boldsymbol{D}(\boldsymbol{u}) \quad \text{in } \Omega$$

< 1

$$oldsymbol{D}(oldsymbol{u}) = rac{1}{2} (
abla oldsymbol{u} + (
abla oldsymbol{u})^T), \quad \dot{\gamma} = \sqrt{2 oldsymbol{D}(oldsymbol{u}) : oldsymbol{D}(oldsymbol{u}))}$$



Rouleaux aggregation

$$\mu_0 = \lim_{\dot{\gamma} \to 0} \mu(\dot{\gamma}) = 0.056 Pa s$$

$$\mu_\infty = \lim_{\dot{\gamma} \to \infty} \mu(\dot{\gamma}) = 0.00345 Pa s$$

$$\lambda = 3.313 s$$

$$n = 0.3568$$



Shear-thinning viscosity Carreau model $\mu(\dot{\gamma}) = \mu_{22}$

$$\frac{\mu(\gamma) - \mu_{\infty}}{\mu_0 - \mu_{\infty}} = \left[1 + (\lambda \dot{\gamma})^2\right]^{(n-1)/2}, \quad n$$

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MATHEMATICAL RESULTS

If the correspondence $T \to \mu(|T|^2)T$ satisfies, for some p > 1, and for every symetric second order tensors A, B, the relations

$$\begin{split} & \mu(|\boldsymbol{A}|^2)\boldsymbol{A}: \boldsymbol{A} \geq C|\boldsymbol{A}|^p \\ & |\mu(|\boldsymbol{A}|^2)\boldsymbol{A}| \leq C(1+|\boldsymbol{A}|)^{p-1} \\ & \left(\mu(|\boldsymbol{A}|^2)\boldsymbol{A} - \mu(|\boldsymbol{B}|^2)\boldsymbol{B}\right): (\boldsymbol{A} - \boldsymbol{B}) > 0 \\ & \text{then for } f \in L^{p'}(\Omega) \text{ and } p > \frac{2d}{d+1}, \text{ there exists a weak solution} \\ & \boldsymbol{u} \in W^{1,p}(\Omega), \end{split}$$

$$\int_{\Omega} \mu(\dot{\gamma}) \boldsymbol{D} \boldsymbol{u} : \boldsymbol{D} \boldsymbol{v} + \int_{\Omega} (\boldsymbol{u} \cdot \nabla \boldsymbol{u}) \cdot \boldsymbol{v} = \int_{\Omega} \boldsymbol{f} \cdot \boldsymbol{v}, \quad \forall \boldsymbol{v} \in \mathcal{V}.$$
 $p > rac{3d+2}{d+2}$ (unsteady, Ladyzhenskaya (1968-70), Lions(1969))

 $p > \frac{3d}{d+2}$ (unsteady, space periodic, Bellout(1994), Málek(1996))

 $p \ge \frac{3d}{d+2}$ (steady, Lions(1969)) $p > \frac{2d}{d+1}$ (steady, Frehse(1997), Ruzicka(1997), Arada and Sequeira(2005))

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MORPHOLOGY OF THE BLOOD VESSELS



The vessel wall is formed by many layers made of tissues with different mechanical characteristics

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Equations for the deformation of the vessel wall

3D nonlinear hyperelasticity (Lagrangian formulation)





 $\partial \Sigma^{t} = \Gamma^{t}_{\omega} \cup \Gamma^{t}_{\Sigma,ext} \cup \Gamma^{t}_{\Sigma,in} \cup \Gamma^{t}_{\Sigma,out}$ $\partial \Sigma^{0} = \Gamma^{0}_{\omega} \cup \Gamma^{0}_{\Sigma,ext} \cup \Gamma^{0}_{\Sigma,in} \cup \Gamma^{0}_{\Sigma,out} \longrightarrow \text{ reference boundaries}$





3D nonlinear hyperelasticity (Lagrangian formulation)

$$\rho_{w} \frac{\partial^{2} \eta}{\partial t^{2}} - \nabla_{0} \cdot \sigma(\eta) = 0 \quad \text{in } \Sigma^{0}, \forall t \in I$$



 $\rho_w \longrightarrow$ wall density



first Piola-Kirchhoff tensor

deformation gradient tensor

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second Piola-Kirchhoff tensor



3D nonlinear hyperelasticity (Lagrangian formulation)

Green-St Venant strain tensor

$$E = E(\eta) = \frac{1}{2} (F^T F - I) = \frac{1}{2} ((\nabla_0 \eta)^T + \nabla_0 \eta + (\nabla_0 \eta)^T \nabla_0 \eta)$$

St Venant – Kirchhoff material

 $S(\eta) = \lambda tr(E)I + 2\nu E$ (linear response)

with







Equations for the deformation of the vessel wall (Lagrangian formulation)

 $\rho_{w} \frac{\partial^{2} \eta}{\partial t^{2}} - \nabla_{0} \cdot \sigma(\eta) = 0 \quad \text{in } \Sigma^{0}$ $\eta = \eta_0$ for t = 0, in Σ^0 + compatibility conditions $\frac{\partial \eta}{\partial t} = \frac{\partial \eta_0}{\partial t} \quad \text{for } t = 0, \quad \text{in } \Sigma^0$ $\frac{\partial \eta_0}{\partial t} = u_0 \quad \text{on } \Gamma^0_{\omega}$ $\sigma(\eta).n_0 = \hat{\phi} \qquad \text{on } \Gamma_{\omega}^0$ & $\sigma(\eta).n_0 = 0 \qquad \text{on } \Gamma^0_{\Sigma.ext}$ boundary conditions interface conditions $\eta = 0$ on $\Gamma^0_{\Sigma,out}$ on $\Gamma^0_{\Sigma,in}$ $\eta = 0$ (clamped structure)

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initial

&



• Blood flow: Generalized Newtonian flow (ALE frame)



Deformation of the vessel wall

$$\rho_{w} \frac{\partial^{2} \eta}{\partial t^{2}} - \nabla_{0} \cdot \sigma(\eta) = 0 \quad \text{in } \Sigma^{0}$$

Interface conditions

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$$oldsymbol{\sigma}(\eta) \cdot oldsymbol{n} = -poldsymbol{n} + oldsymbol{ au}(oldsymbol{u}) \cdot oldsymbol{n}$$
 at Γ_w
 $oldsymbol{u} = rac{\partial \eta}{\partial t}$ at Γ_w



u = blood velocityw = domain velocity p = pressure $\rho_f = density$ $\mu = viscosity$ $\eta = wall displacement$

+ initial and boundary conditions at Γ_i (i=0,1,2)



Interface conditions

$$u = \frac{\partial \eta}{\partial t}, \quad \forall t \in I, \text{ at } \Gamma_{\omega}^{t}$$

$$\sigma(\eta) \cdot n = -pn + \tau(u) \cdot n, \quad \forall t \in I, \text{ at } \Gamma_{\omega}^{t}$$



(using the Piola transform)

 $-(\det \nabla_0 \eta)\tau(u,p)(\nabla_0^{-T}\eta) \cdot n_0 = \sigma(\eta) \cdot n_0, \quad \forall t \in I, \text{ on } \Gamma_{\omega}^t$





Blood:

Newtonian or non-Newtonian fluid

Deformation of the Vessel Wall:

3D (nonlinear) elasticity or 2D shell type models



Displacement (new domain)



Normal stress

implicit coupling (iterative procedure)



Open problems:

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Well posedness of the FSI problem

Contributions given by e.g. : D.Coutand, S. Shkoller, Y.Maday, C.Grandmont, B.Desjardins, M.Esteban, G.P. Galdi, H.Beirão da Veiga, among others

Devise efficient numerical algorithms

Contributions given by e.g. : P. le Tallec, F.Nobile, M.A.Fernandéz,M.Moubachir,J-F.Gerbeau, S.Deparis, W.A.Wall, among others



Regularity Assumptions:

 $\Omega^t \subset \mathbb{R}^3$ is open and connected

 $\partial \Omega^t = \Gamma^t_\omega \cup \Gamma^t_{in} \cup \Gamma^t_{out}$ Is locally Lipschitz ($\partial \Omega^t \in C^{1,1}$)

$$\frac{\partial \eta}{\partial t} \in H^{1/2}(\Gamma_{\omega}^{t}) \implies u(t) \in H^{1}(\Omega^{t}), \forall t$$





An Energy Estimate for the coupled FSI problem

[J. Janela, A. Moura, A. S, 2009 – generalization of L. Formaggia, A. Moura, F. Nobile, 2007]

$$\mathbf{E}(t) = \frac{\rho}{2} \|u\|_{L^{2}(\Omega^{t})}^{2} + \frac{\rho_{w}}{2} \left\|\frac{\partial\eta}{\partial t}\right\|_{L^{2}(\Sigma^{0})}^{2} + \mu(\gamma) \|E(\eta)\|_{L^{2}(\Sigma^{0})}^{2} + \frac{\lambda}{2} \|trE(\eta)\|_{L^{2}(\Sigma^{0})}^{2}$$

THEOREM: The coupled FSI problem, with homogeneous Dirichlet boundary conditions u = 0 at Γ_{in}^{t} and Γ_{out}^{t} satisfies the following energy inequality

$$\frac{d}{dt}(\mathbf{E}(t)) + 2\mu_{\infty} \left\| D(u) \right\|_{L^{2}(\Omega^{t})}^{2} \le 0 \quad \text{and, consequently, the}$$

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 $E(t) + 2\mu_{\infty} \int_{0}^{\infty} \|D(u)\|_{L^{2}(\Omega')}^{2} dt \leq E(0)$ where E(0) is a constant depending only on the initial data $u_{0}, \eta_{0}, \eta_{0}$

REMARK: $\int_{\Gamma'_{in}} |u|^2 u.n > 0$, $\int_{\Gamma'_{out}} |u|^2 u.n > 0$ for homogeneous Neumann conditions



energy decay property

Sketch of the PROOF:

- 1. Multiply the structure equation by $\frac{\partial \eta}{\partial t}$, integrate over the reference domain, use the boundary and matching conditions
- 2. Multiply the fluid equation by u, integrate over the fluid domain, ...

3.
$$\mu_{\infty} \le \mu(\gamma) \le \mu_0$$
 (shear-thinning viscosity fluid)

$$= \int_{\Omega^{t}} 2\mu(\gamma)D(u): \nabla u \, d\omega \ge 2\mu_{\infty} \left\| D(u) \right\|_{L^{2}(\Omega^{t})}^{2}$$

$$= \frac{d}{dt}(\mathbf{E}(t)) + 2\mu_{\infty} \left\| D(u) \right\|_{L^{2}(\Omega^{t})}^{2} + \frac{\rho}{2} \int_{\mathbb{T}^{t}} |u|^{2} u.nd\gamma \le \int_{\mathbb{T}^{t}} (\tau(u).n).ud\gamma$$

 $\Gamma_{in}^{i} \cup \Gamma_{out}^{i}$

 $\Gamma_{in}^{*} \cup \Gamma_{out}^{*}$

Finally

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FSI Algorithm: (adapted from Fernandéz & Moubachir, 2005)

ALE formulation to account for the evolution of the computational domain

Efficient solvers for each fluid and structure subproblems to ensure accurate and fast convergence of the FSI nonlinear coupled system

Fluid equations:	Discretization in time: implicit Euler scheme	
	Discretization in space: stabilized P1 bubble / P1 FE	
Structure equations:	Discretization in time: mid-point Newmark method	
	Discretization in space: P1 FE	

Coupling strategy: fully implicit coupling based on a Newton algorithm with the exact computation of the Jacobian





Newtonian vs non-Newtonian behavior

Main objectives:

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- 3D Non-Newtonian models for blood flow
- 3D Fluid-Structure Interaction algorithms for pressure wave propagation in arteries and detailed flow patterns using Newtonian and non-Newtonian blood flow models
- Geometrical multiscale simulation of the cardiovascular system using non-Newtonian models



Idealized vs reconstructed geometries & computational grids



carotid bifurcation







Newtonian vs non-Newtonian

Carotid Bifurcation: Wall Shear Stress (WSS)



Carreau model Curved vessel: Pressure



Structure:

 $\lambda(E,\xi) = 3x10^6 \, dynes \, / \, cm^2$ $\nu(E,\xi) = 0.3$ $\rho_{\omega} = 1.2g \, / \, cm^2$ h = 0.1cm

A. Moura & J. Janela



Fluid:

 $\mu_{0} = \lim_{\substack{\gamma \to 0 \\ \gamma \to 0}} \mu(\gamma) = 0.056 Pas$ $\mu_{\infty} = \lim_{\substack{\gamma \to \infty \\ \gamma \to \infty}} \mu(\gamma) = 0.00345 Pas$ $\lambda = 3.313 s$ n = 0.3568 $\rho = 1g / cm^{3}$

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Newtonian vs non-Newtonian Curved vessel: Wall Shear Stress (WSS)



Carotid Bifurcation: Pressure pulse



Carreau model

A. Moura & J. Janela





Newtonian vs non-Newtonian

Carotid Bifurcation: Wall Shear Stress (WSS)



GEOMETRICAL MULTISCALE



•Global features have influence on the local fluid dynamics

 Local changes in geometry or material properties (e.g. due to surgery, aging, stenosis, …) may induce pressure waves reflections
 → global effects

Modeling strategy

- use the expensive 3D model only in the region of interest
- couple with network models that include peripheral impedances to account for global effects



GEOMETRICAL MULTISCALE



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3D

- Very detailed simulations
- Very complex
- Computationally very costly

1D

- Evolution of mean pressure and flux in arteries
- System of hyperbolic equations
- Low computational cost

- Evolution in time of mean pressure and flux in wide compartments
- System of ODEs
- Very low computational cost



GEOMETRICAL MULTISCALE 1D Model

M

Allows for the simulation of complex arterial networks!

Domain decomposition



$$\begin{cases} \frac{\partial A}{\partial t} + \frac{\partial Q}{\partial z} = 0\\ \frac{\partial Q}{\partial t} + \alpha \frac{\partial}{\partial z} \left(\frac{Q^2}{A}\right) + \frac{A}{\rho} \frac{\partial P}{\partial z} + K \frac{Q}{A} = 0\\ P - P_0 = \Psi(A) \end{cases}$$
Area $\longrightarrow A(z,t) = \int_{\Omega \cap \Sigma(z)} d\gamma$
Flux $\longrightarrow Q(z,t) = \int_{\Omega \cap \Sigma(z)} u_z(x,t) d\gamma$
can Pressure $\longrightarrow P(z,t) = \frac{1}{|\Sigma(z)|} \int_{\Omega \cap \Sigma(z)} p(x,t) d\gamma$

describes de wave propagation nature of blood flow

- acts as absorbing boundary condition for the 3D model
- simulation of complex arterial trees by coupling 1D models





GEOMETRICAL MULTISCALE 0D Model



CCINAT IST CENT MATEMA E APLICA 0D Lumped parameters (system of linear ODE's)

$$C\frac{dP_i}{dt} = -(Q_{i+1} - Q_i),$$

$$L\frac{dQ_i}{dt} = -(P_i - P_{i-1}) - RQ_i$$

The analogy	Fluid dynamics	Electrical circuits	
	Pressure	Voltage	
	Flow rate	Current	
	Blood viscosity	Resistance R	
	Blood inertia	Inductance L	
	Wall compliance	Capacitance C	

• RLC circuits model "large" arteries

- RC circuits account for capillary bed
- Can describe compartments (such as peripheral circulation)



GEOMETRICAL MULTISCALE

3D and 1D for a cylindrical artery: pressure pulse



GEOMETRICAL MULTISCALE

3D-1D for the carotid bifurcation: velocity field & pressure pulse



(A. Moura)





MODELING CEREBRAL ANEURYSMS

Cerebral Aneurysms:

- Most common cause of hemorrhagic strokes
- Tend to be silent until rupture
- High prevalence, low risk

Main Goal:

 Help improve the evaluation & treatment of cerebral aneurysms

Our Approach:

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 Patient-specific image-based CFD modeling to link hemodynamics & clinical observations







MECHANISMS

- The mechanisms responsible for the development, growth and rupture of intracranial aneurysms are not well understood
- Better understanding of these processes can lead to better patient evaluation and improved treatments







IMAGE-BASED MODELING OF BLOOD FLOWS







FLOW COMPLEXITY & STABILITY

simple



J. Cebral, George Mason Univ

complex





DATABASE: ANEURYSM MODELS & CLINICAL INFO



J. Cebral, George Mason Univ

CGI

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CIRCLE of WILLIS: 1D – NETWORK ?







CIRCLE of WILLIS: 1D - NETWORK



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MRA-BASED SUBJECT- SPECIFIC MODELING



FINITE ELEMENT MESH

Advancing front method >20 million tetrahedra







CONCLUSIONS/ OUTCOME

Patient-specific CFD models are capable of realistically representing the *in vivo* hemodynamic characteristics

These models can be used to better understand the mechanisms of aneurysm growth and rupture

They can also be used to answer specific clinical questions and to improve aneurysm risk assessment

Simulation-assisted treatment planning and patient evaluation tools are becoming a reality





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CURRENT PROJECTS

Multiscale Mathematical Modelling in Biomedicine PTDC/ MAT / 68166/ 2006 [2007–2010]

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 Cardiovascular Imaging Modeling and Simulation – SIMCARD UTAustin/CA/0047/2008 [2009 - 2012]



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THANK YOU !!!



